Marko Berghoff, Humboldt Universität zu Berlin

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## Introduction

QFT in position space / Causal perturbation theory

 Stueckelberg, Bogoliubov, Shirkov (late 50's): Axiomatic approach to S-matrix,

$$S=1+\sum_{n>1}T_n.$$

- Epstein and Glaser ('73): Renormalization of S translates into an extension (splitting) problem for distributions.
- Simplified version by Stora (ca.'00), used in QFT on curved spacetimes.

## Introduction

Bergbauer, Brunetti, Kreimer ('10): Version for single graphs.

Example (Euclidean  $\phi_4^4$ -theory)



Feynman rules  $\Phi: G \mapsto \int \omega_G = \int dx dy dz \frac{1}{(x-y)^4 (y-z)^2 z^4 x^2}$ What is  $\int \omega_G$ ?

- Easy answer:  $\infty$ .
- Tricky answer: Find renormalized value ...

Idea (Atiyah; Axelrod, Singer): Use a smooth model to arrange the divergences in a "nice" way, renormalize on this model, then push the result back to original spacetime.

### Definition

Let  $\mathcal{A} = \{A_1, \ldots, A_k\}$  be a family of smooth subvarities in an algebraic variety X. A smooth model is a smooth variety Y together with a proper, surjective map  $\beta : Y \to X$ , such that  $\mathcal{E} := \beta^{-1}(\cup_{A \in \mathcal{A}} A)$  is a normal crossing divisor and  $\beta_{|Y \setminus \mathcal{E}}$  a diffeomorphism.

Such smooth models are given by the wonderful model construction by DeConcini and Procesi. Idea is based on Fulton and MacPherson's "Compactification of Configuration Spaces":

The configuration space of n-points in an algebraic variety X is

$$C_n(X) = \{ (x_1, \ldots, x_n) \in X^n \mid x_i \neq x_j \text{ for all } i \neq j \}.$$

Fulton and MacPherson construct its compactification X[n] by a sequence of blow-ups along the (strict transforms) of diagonals of increasing dimension. A limiting point in  $X[n] \setminus C_n(X)$  is encoded by a nested set of diagonals.

# Wonderful models

#### Definition

Let  $\mathcal{A}$  be a linear arrangement in a vector space X. The wonderful model  $(Y_{\mathcal{A}}, \beta)$  is defined as follows: The graph of the map

$$\pi_{\mathcal{A}}: X \setminus \bigcup_{A \in \mathcal{A}} A \longrightarrow \prod_{A \in \mathcal{A}} \mathbb{P}(X/A)$$

is locally closed in  $X \times \prod_{A \in \mathcal{A}} \mathbb{P}(X/A)$ . Define  $Y_{\mathcal{A}}$  as its closure and  $\beta : Y_{\mathcal{A}} \to X$  as the projection onto the first factor.

- ▶ An explicit construction is given by a sequence of blowups along (strict transforms of) elements of a building set  $\mathcal{B} \subseteq \mathcal{A}$ , giving local charts  $(U_{\underline{i}}, \kappa_{\underline{i}})$ ,  $\underline{i} = (\mathcal{N}, B)$ , where  $\mathcal{N}$  is a nested set of elements of  $\mathcal{B}$  and B an adapted, marked basis of X.
- B controls the number of irreducible components of E ⊆ Y<sub>B</sub>, while the B-nested sets describe a stratification of E.

## Graphs and arrangements

Feichtner ('05): These notions can all be defined combinatorially! Either in terms of the intersection lattice of A,

$$\mathcal{L}_{\mathcal{A}} := \Big\{ \{A_1 \cap \cdots \cap A_k \mid A_i \in \mathcal{A}\}, \supseteq \Big\},$$

or, in our case, using the poset of divergent subgraphs of G.

#### Definition

Let G = (V, E) be a graph.

- Its superficial degree of divergence is defined by
  s(G) = dh<sub>1</sub>(G) 2|E| (d = dim. of spacetime). G is called at most logarithmic if s(g) ≤ 0 holds for all g ⊆ G.
- The divergent poset of G is defined as

$$\mathcal{D}_{G} := \Big\{ \{ g \subseteq G \mid s(g) \leq 0 \}, \subseteq \Big\}.$$

Now consider the following Feynman rules:

Let G be a connected graph. Orient G and choose a spanning tree  $t \subseteq G$ .

The Feynman rules map  $\Phi$  sends G to the pair  $(X_G, \tilde{v}_G)$  of a chain  $X_G = (\mathbb{R}^d)^{E(t)}$  and a form defined by the rational function

$$v_G: x \longmapsto \prod_{e \in E(G)} y_e^{-\frac{d}{2}}, \quad y_e = \begin{cases} x_e & \text{if } e \in E(t) \\ \sum_{e' \in E(t_e)} \sigma_t(e') x_{e'} & \text{else.} \end{cases}$$

Here  $t_e$  is the unique path in t connecting the source and target vertices of e and  $\sigma_t : E(t) \to \{\pm 1\}$  given by the orientation on G.

# Graphs and arrangements

We avoid the infrared problem -  $v_G \notin L^1(X_G)$  - by viewing  $v_G$  as (the kernel of) a distribution on  $X_G$ . On the other hand, the ultraviolet problem -  $v_G \notin L^1_{loc}(X_G)$  - is characterized by the following

### Proposition

Let G be at most logarithmic.

- ▶  $v_G$  defines a distribution on  $X_G \setminus \bigcup_{g \in D_G} A_g$ , where  $A_g := \{y_e = 0 \mid e \in E(g)\} \subseteq X_G$ .
- ▶ D<sub>G</sub> is a graded (distributive) lattice with join and meet operations given by

$$g \lor h := g \cup h$$
$$g \land h := g \cap h.$$

# Wonderful combinatorics

### Definition

Let  $\mathcal{L}$  be a lattice.  $\mathcal{B} \subseteq \mathcal{L}$  is a building set for  $\mathcal{L}$  if

▶ for all  $p \in \mathcal{L}_{>\hat{0}}$  and  $\{q_1, \ldots, q_k\} = \max \mathcal{B}_{\leq p}$  there is an isomorphism of posets

$$\varphi_{A}: \prod_{i=1}^{k} [\hat{0}, q_{i}] \longrightarrow [\hat{0}, p]$$

with 
$$\varphi_p(\hat{0},\ldots,q_j,\ldots,\hat{0}) = q_j$$
 for  $j = 1,\ldots,k$ .

▶ the ranking function on *L* satisfies

$$r(p) = \sum_{i=1}^k r(q_i).$$

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In our case r is given by  $\operatorname{codim}(A_g) = d(|E(g)| - h_1(g))$ .

# Wonderful combinatorics

#### Definition

Let  $\mathcal{B}$  be a building set in  $\mathcal{L}$ . A subset  $\mathcal{N} \subseteq \mathcal{B}$  is  $\mathcal{B}$ -nested if for all subsets  $\{p_1, \ldots, p_k\} \subseteq \mathcal{N}$  of pairwise incomparable elements their join  $p_1 \vee \cdots \vee p_k$  exists (in  $\mathcal{L}$ ) and does not belong to  $\mathcal{B}$ .

### Definition

- A basis b of X is adapted to N if for all A ∈ N the set b ∩ A generates A ⇔ b is given by the edges of an adapted spanning tree, i.e. t ⊆ G such that t ∩ g is spanning for all g ∈ N.
- A marking of b is for every A ∈ N the choice of an element b<sub>A</sub> ∈ b ∩ A ⇐⇒ for every g ∈ N a choice b<sub>g</sub> ∈ {b<sub>e</sub>}<sub>e∈E(t∩g)</sub>.

Let  $(Y_{\mathcal{B}}, \beta)$  be a wonderful model for a building set  $\mathcal{B} \subseteq \mathcal{D} = \mathcal{D}_G$ and  $v = v_G$  the Feynman distribution associated to a graph G.

#### Proposition

In local coordinates on  $U_{\underline{i}}, \underline{i} = (\mathcal{N}, B)$ , the pullback of  $\tilde{v}^s = v^s |dx|$ (s = regularization parameter) along  $\beta$  is a density on Y given by

$$\begin{split} (\omega^{s})_{\underline{i}} &:= (\beta^{*}\tilde{v}^{s})_{\underline{i}} = f_{\underline{i}}^{s} \prod_{g \in \mathcal{N}} u_{g}(s, \cdot) |dx|, \\ u_{g}(s, x) &= |x_{g}|^{-1 + r(g)(s-1)}, x_{g} \text{ marked}. \end{split}$$

The map  $f_{\underline{i}} : \kappa_{\underline{i}}(U_{\underline{i}}) \longrightarrow \mathbb{R}$  is in  $L^1_{loc}(\kappa_{\underline{i}}(U_{\underline{i}}))$  and smooth in the marked variables  $x_g$ ,  $g \in \mathcal{N}$ .

The next step is to study the Laurent expansion of  $\omega^s$ . To formulate this we need a local version of graph contraction.

#### Definition

Let  $g \subseteq G$  and  $\mathcal{N}$  be nested. The contraction relative to  $\mathcal{N}$  is defined as

$$g/\!/\mathcal{N} := egin{cases} g/(igcup_{\gamma\in\mathcal{N}_{< g}}\gamma) & ext{if } g\in\mathcal{N}, \ g/(g\capigcup_{\gamma\in\mathcal{N},\gamma\cap g< g}\gamma) & ext{else.} \end{cases}$$

For  $\mathcal{J}\subseteq\mathcal{N}$  the poset  $(\mathcal{N}/\!/\mathcal{J},\sqsubseteq)$  is given by the underlying set

$$\mathcal{N}//\mathcal{J} := \{g//\mathcal{J} \mid g \in \mathcal{N}\},$$

partially ordered by inclusion (in general  $\sqsubseteq \neq \subseteq !$ ).

## Theorem

- The Laurent expansion of ω<sup>s</sup> at s = 1 has a pole of order N where N is the cardinality of the largest B-nested set.
- ► The coefficients ã<sub>k</sub> in the principal part of the Laurent expansion

$$\omega^{s} = \sum_{-N \leq k} \tilde{a}_{k} (s-1)^{k}$$

are densities with supp  $\tilde{a}_k = \bigcup_{|\mathcal{N}|=-k} \mathcal{E}_{\mathcal{N}}$  for k < 0.

Consider the minimal building set I(D) ⊆ D. Assume G ∈ I(D). Let N be the cardinality of a maximal nested set and denote by χ the constant function on the wonderful model Y<sub>I(D)</sub>. Then

$$\langle \tilde{a}_{-N} | \chi \rangle = \sum_{|\mathcal{M}|=N} \prod_{\gamma \in \mathcal{M}} \mathcal{P}(\gamma / / \mathcal{M}).$$

### Definition ("Local subtraction at fixed conditions")

In every chart  $U_{\underline{i}}$  let  $\nu = \{\nu_{\underline{g}}^{\underline{i}}\}_{g \in \mathcal{N}}$  denote a collection of smooth functions on  $\kappa(U_{\underline{i}})$ , each  $\nu_{\overline{g}}^{\underline{i}}$  depending only on the coordinates  $x_e$  with  $e \in E(t) \cap E(g \setminus \mathcal{N}_{\leq g})$ , satisfying  $\nu_{\overline{g}}^{\underline{i}}|_{x_g=0} = 1$  and compactly supported in all other directions. For  $u \in \mathcal{D}'(\mathbb{R} \setminus \{0\})$  and  $\mu \in \mathcal{D}([-1,1])$  let  $r_{\mu}[u] \in \mathcal{D}'(\mathbb{R})$  denote the extended distribution

$$r_{\mu}[u]: \varphi \mapsto \langle u|\varphi \rangle - \langle u|\varphi(0)\mu \rangle.$$

The extension of  $\omega^{s}$  is defined by

$$\begin{aligned} \mathcal{R}_{\nu}[\omega^{s}] \stackrel{\textit{loc.}}{=} & \mathcal{R}_{\nu}^{i}[f_{\underline{i}}^{s} \prod_{g \in \mathcal{N}} u_{g}(s) | dx |] := f_{\underline{i}}^{s} \prod_{g \in \mathcal{N}} r_{\nu_{g}^{i}}[u_{g}(s)] | dx | \\ &=: \sum_{\mathcal{J} \subseteq \mathcal{N}} (-1)^{|\mathcal{J}|} \nu_{\mathcal{J}}^{i} \cdot (\omega_{\underline{i}}^{s})_{\mathcal{E}_{\mathcal{J}}}. \end{aligned}$$

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#### Theorem

- R<sub>ν</sub>[ω<sup>s</sup>] defines a density-valued holomorphic function in a neighborhood of s = 1.
- Define the renormalized Feynman rules by the map

$$\Phi_R: G \longmapsto (X_G, \mathscr{R}[v_G])$$

with  $\mathscr{R}[v_G] := \beta_* R_{\nu}[\omega^s]_{|s=1}$  and evaluation on  $\varphi \in \mathcal{D}(X_G)$  given by

 $\langle \mathscr{R}[\mathsf{v}_G] \mid \varphi \rangle = \langle \beta_* R_{\nu}[\omega^s]_{|s=1} \mid \varphi \rangle = \langle R_{\nu}[\omega^s]_{|s=1} \mid \beta^* \varphi \rangle.$ 

Then  $\mathscr{R}$  satisfies the Epstein-Glaser locality principle.

## Renormalization group

What happens if the renormalization point  $\nu$  is changed?

Theorem

Consider  $(R_{\nu'} - R_{\nu})[\omega^s]$  for two choices of function families  $\nu'$  and  $\nu$ . Locally in  $U_{\underline{i}}$ , applied on a test function  $\varphi = \beta^* \psi$  for  $\psi \in \mathcal{D}(\beta(U_{\underline{i}}) \cap \kappa_{\underline{i}}(U_{\underline{i}}))$  we have

$$\langle (R^{\underline{i}}_{\nu'} - R^{\underline{i}}_{\overline{\nu}})[\omega^{s}_{\underline{i}}]|\varphi\rangle = \sum_{\emptyset \neq \mathcal{J} \subseteq \mathcal{N}} c_{\mathcal{J}} \langle R^{\underline{j}}_{\overline{\nu}}[(\omega^{s}_{G/\!/\mathcal{J}})_{\underline{j}}]|\delta_{\mathcal{J}}[\varphi]\rangle$$

with

$$c_{\mathcal{J}} := \prod_{\gamma \in \mathcal{J}} \langle \mathsf{R}^{\underline{k}}_{\nu} [(\omega^{s}_{\gamma / / \mathcal{J}})_{\underline{j}}] \mid \nu_{\gamma}' \rangle.$$

The indices  $\underline{j}, \underline{k}$  correspond to  $(\mathcal{N}//\mathcal{J})_{\sqsubseteq G//\mathcal{J}}$  and  $(\mathcal{N}//\mathcal{J})_{\sqsubseteq \gamma//\mathcal{J}}$ , respectively.

- Geometric ansatz put in combinatorial language
- Simplifies the "wonderful" construction and adds discrete toolbox
- ► Reconstruction of Epstein-Glaser method via models for K<sub>n</sub> → Fulton-MacPherson compactification
- Dyson-Schwinger equations?
- Renormalization group equation / flow?
- Renormalization Hopf algebra? It encodes the stratification of the exceptional divisor *E*...

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